

Non Hermetian Quantum Dynamics of Damped Oscillator

Kushagra Nigam

University of Kentucky

kushagran1@gmail.com

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THE BORDER TERRITORY

QUANTUM DOMAIN

CLASSICAL DOMAIN



Classical Damped Harmonic Oscillator

- **Classical Equation**

$$\ddot{q} + 2\gamma\dot{q} + \omega^2q = 0 \quad (1)$$

where,

\dot{q} indicates time derivative

γ is damping coefficient

ω is oscillator frequency

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- **Solution**

$$q(t) = q_o e^{-\gamma t} \cos \omega_d t + \left(\frac{\dot{q}_o}{\omega_d} + \frac{\gamma q_o}{\omega_d} \right) e^{-\gamma t} \sin \omega_d t \quad (2)$$

where, $\omega_d = \sqrt{\omega^2 - \gamma^2}$ and assume $p_o = m\dot{q}_o + m\gamma q_o$

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- **Quantum Average using Ehrenfest theorem**

$$\langle \alpha, t | \hat{q} | \alpha, t \rangle = e^{-\gamma t} \left[\langle \alpha | \hat{q} | \alpha \rangle \cos \omega_d t + \frac{\langle \alpha | \hat{p} | \alpha \rangle}{m\omega_d} \sin \omega_d t \right] \quad (3)$$

where, $|\alpha, t\rangle$ is coherent state

- **General Non-Hermetian Hamiltonian**

$$\hat{H}_D = \hat{H} - i\hat{\Gamma} \quad (4)$$

where,

$\hat{H} = \hat{H}^\dagger$ is Hamiltonian of undamped system

$\hat{\Gamma} = \hat{\Gamma}^\dagger$ is damping part of Hamiltonian

$$[\hat{H}, \hat{\Gamma}] = 0$$

- **Time Evolution**

$$i\hbar \frac{\partial |\psi\rangle}{\partial t} = \hat{H}_D |\psi\rangle \quad (5)$$

$$\implies |\psi, t + dt\rangle = \hat{U}(t, t + dt) |\psi, t\rangle = \left(1 - i \frac{\hat{H}_D dt}{\hbar} \right) |\psi, t\rangle \quad (6)$$

where,

$\hat{U}(t, t + dt)$ is non-unitary time evolution operator

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Similarly,

$$|\psi, t\rangle = \hat{U}(t) |\psi, 0\rangle = e^{-i \frac{\hat{H}_D t}{\hbar}} |\psi, 0\rangle \quad (7)$$

- **Consider a state of damped system**

$$|\psi, t = 0\rangle = \sum_{n=0}^{\infty} c_n |n\rangle \quad (8)$$

where, $|n\rangle$ are energy eigen states of undamped Hamiltonian and

$$\sum_{n=0}^{\infty} |c_n|^2 = 1$$

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$$|\psi, t\rangle = \hat{U}(t) |\psi, 0\rangle = e^{-i\frac{\hat{H}_D t}{\hbar}} |\psi, 0\rangle \quad (9)$$

$$= \sum_{n=0}^{\infty} c_n e^{-\frac{\Gamma_n t}{\hbar}} e^{-i\frac{\hat{E}_n t}{\hbar}} |n\rangle \quad (10)$$

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- Time Dependent Norm!!

$$N(t) = \langle \psi, t | \psi, t \rangle = \sum_{n=0}^{\infty} |c_n|^2 e^{-2\frac{\Gamma_n t}{\hbar}} \quad (11)$$

- Time Derivative of Norm

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- Operator Expectation Values

$$\langle \hat{A} \rangle_t = \frac{\langle \psi, t | \hat{A} | \psi, t \rangle}{N(t)} \quad (13)$$

$$= \langle \phi, t | \hat{A} | \phi, t \rangle \quad (14)$$

where, $|\phi, t\rangle = \frac{|\psi, t\rangle}{\sqrt{N(t)}}$

- Evolution of Expectation Values

$$\langle \psi, t + dt | \hat{A} | \psi, t + dt \rangle = \langle \psi, t | \hat{U}^\dagger(dt) \hat{A} \hat{U}(dt) | \psi, t \rangle \quad (15)$$

$$\implies i\hbar \frac{d \langle \psi, t | \hat{A} | \psi, t \rangle}{dt} = \langle \psi, t | [\hat{H}, \hat{A}] | \psi, t \rangle - i \langle \psi, t | \{ \hat{\Gamma}, \hat{A} \} | \psi, t \rangle \quad (16)$$

where, $[\hat{H}, \hat{A}] = \hat{H}\hat{A} - \hat{A}\hat{H}$ and $\{ \hat{\Gamma}, \hat{A} \} = \hat{\Gamma}\hat{A} + \hat{A}\hat{\Gamma}$

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But expectations have to be calculated using normalized states, hence

$$i\hbar \frac{d \langle \phi, t | \hat{A} | \phi, t \rangle}{dt} = \langle \phi, t | [\hat{H}, \hat{A}] | \phi, t \rangle - 2i\Delta_{A\Gamma}^2 \quad (17)$$

where, $\Delta_{A\Gamma}^2 = \langle \phi, t | \frac{1}{2} \{ \hat{A}, \hat{\Gamma} \} | \phi, t \rangle - \langle \phi, t | \hat{A} | \phi, t \rangle \langle \phi, t | \hat{\Gamma} | \phi, t \rangle$

Modified Schrödinger Equation

$$i\hbar \frac{d|\phi, t\rangle}{dt} = \hat{H}|\phi, t\rangle - i[\hat{\Gamma}, |\phi, t\rangle \langle \phi, t|]|\phi, t\rangle \quad (18)$$

where, $|\phi, t\rangle$ are normalized kets

Coherent States

- For our case, we choose

$$\hat{\Gamma} = \hbar\gamma\hat{N} \quad (19)$$

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$$|\alpha, 0\rangle = \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} |n\rangle \quad (21)$$

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- Non Unitary Time Evolution

$$|\alpha, t\rangle = \hat{U}(t) |\alpha, 0\rangle = e^{-i\frac{\hat{H}_D t}{\hbar}} |\alpha, 0\rangle = \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} e^{-\gamma n t} e^{-i\frac{\hat{E}_n t}{\hbar}} |n\rangle \quad (22)$$

- Time Dependent Norm

$$N(t) = \langle \alpha, t | \alpha, t \rangle = e^{|\alpha|^2} e^{-2\gamma t} \quad (23)$$

Normalized Coherent States

$$|\alpha, t\rangle = e^{-\frac{|\alpha|^2}{2}} e^{-2\gamma t} \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} e^{-\gamma n t} e^{-i\frac{\hat{E}_n t}{\hbar}} |n\rangle \quad (24)$$

- Position Expectation

$$\begin{aligned} q(t) &= \langle \alpha, t | \hat{q} | \alpha, t \rangle = \langle \alpha, t | \sqrt{\frac{\hbar}{2m\omega_d}} (\hat{a} + \hat{a}^\dagger) | \alpha, t \rangle \\ &= \sqrt{\frac{2\hbar}{m\omega_d}} \operatorname{Re}(\alpha e^{-it\omega_d}) e^{-\gamma t} \end{aligned} \quad (25)$$

Note that, $\operatorname{Re}(\alpha) = \sqrt{\frac{m\omega_d}{2\hbar}} q_0$

Quantum to Classical Transition

• Position Expectation

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• Momentum Expectation

$$\begin{aligned} p(t) &= \langle \alpha, t | \hat{p} | \alpha, t \rangle = \langle \alpha, t | \sqrt{\frac{m\omega_d \hbar}{2}} \left(\frac{\hat{a} - \hat{a}^\dagger}{i} \right) | \alpha, t \rangle \\ &= \sqrt{2m\omega_d \hbar} \operatorname{Im}(\alpha e^{-it\omega_d}) e^{-\gamma t} \end{aligned} \quad (26)$$

Note that, $\operatorname{Im}(\alpha) = \sqrt{\frac{1}{2m\omega_d \hbar}} p_0$

Quantum to Classical Transition

- **Position Variance**

$$\begin{aligned}\langle \Delta \hat{q}^2 \rangle &= \frac{\hbar}{2m\omega_d} + \frac{\hbar}{m\omega_d} \left[|\alpha|^2 + \operatorname{Re}(\alpha^2 e^{-it2\omega_d}) - 2 (\operatorname{Re}(\alpha e^{-it\omega_d}))^2 \right] e^{-2\gamma t} \\ &= \frac{\hbar}{2m\omega_d}\end{aligned}\tag{27}$$

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Therefore, $|\alpha, t\rangle$ form minimum uncertainty states.

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Therefore, $|\alpha, t\rangle$ form minimum uncertainty states.

- Heisenberg's Uncertainty is Safe!

$$\langle \alpha, t | \Delta \hat{x}^2 | \alpha, t \rangle \langle \alpha, t | \Delta \hat{p}^2 | \alpha, t \rangle = \frac{\hbar^2}{4}\quad (29)$$

Quantum to Classical Transition

- **Energy Expectation**

$$\begin{aligned}\langle \alpha, t | \hat{H} | \alpha, t \rangle &= e^{-|\alpha|^2} e^{-2\gamma t} \sum_{n,m=0}^{\infty} \frac{|\alpha|^{2n}}{n!} e^{-2\gamma t n} \left[n + \frac{1}{2} \right] \hbar \omega_d \\ &= |\alpha|^2 e^{-2\gamma t} \hbar \omega_d + \frac{1}{2} \hbar \omega_d\end{aligned}\quad (30)$$

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$$i\hbar \frac{d\langle \hat{q} \rangle}{dt} = \langle [\hat{H}, \hat{q}] \rangle - 2i\Delta_{q\Gamma}^2 \quad (31)$$

$$\implies \frac{dq(t)}{dt} = \frac{p(t)}{m} - \gamma q(t) \quad (32)$$

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Differentiate (35) wrt time and use (36) to get,

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but, $\omega_d^2 + \gamma^2 = \omega^2$

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Hence,

Back to Classical Equation

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- **Challenges**

- ① Reservoir Coupling
- ② Measure of Entropy

Danke!